LA-UR-80-2216

# MASTER

TITLE: INTERACTION OF FAR-INFRARED AND MID-INFRARED LASER

TRANSITIONS IN THE AMMONIA LASER

AUTHOR(S): L. Y. Nelson, MSNW

> M. I. Buchwald C. R. Jones

SUBMITTED TO: TOPICAL MEETING ON INFRARED LASERS

- DISCLAIMER .

District on the model of an arrival of seet agreement by an agent, of the binter states (common and in the Binter States (common and an arrival of seet agreement of any agents) (the set agreement) are also given better the arrival of the states (see a section of the second of the s

By acceptance of this article, the publisher recognizes that the U.S. Government rotains a nonexclusive, royalty free license to publish or reproduce the published form of this contribution, or to allow others to do so, for U.S. Government pur

The Los Alamos Scientific Laboratory requests that the publisher identify this article as work performed under the aus pleas of the U.S. Dupartment of Energy

DISTRIBUTION OF THIS DUSDINGHOUTS GREENITED



Post Office Box 1663 Los Alamos, New Mexico 87545 An Affirmative Action/Equal Opportunity Employer

orm No. 836 R3

University of California

# INTERACTION OF FAR-INFRARED AND MID-INFRARED LASER TRANSITIONS IN THE AMMONIA LASER

### L.Y. Nelson

Mathematical Sciences Northwest, Inc. Bellevue, Washington 98004

M.I. Buchwald and C.R. Jones Los Alamos Scientific Laboratory Los Alamos, New Mexico 87545

## **ABSTRACT**

Mid-infrared laser emission in ammonia is usually observed on a P(J+2) transition when a  $CO_2$  laser is used to optically pump a near resonant R(J) absorption feature. However, by generating simultaneous FIR ammonia laser emission in the same optical cavity, mid-infrared emission is obtained exclusively on the P(J) transition.

# INTERACTION OF FAR-INFRARED AND MID-INFRARED LASER TRANSITIONS IN THE AMMONIA LASER

#### L.Y. Nelson

Mathematical Sciences Northwest, Inc. Bellevue, Washington 98004

M.I. Buchwald and C.R. Jones Los Alamos Scientific Laboratory Los Alamos, New Mexico 87545

A new pattern of mid-infrared laser emission has been observed in the optically pumped NH3 laser. Laser emission on the conventional 12.8  $\mu$ m NH3 laser transition, optically pumped by the 9R16 CO2 line, has been completely quenched and replaced by emission at 12.2  $\mu$ m.

Figure 1 presents an energy level diagram for NH3 and the relevant rotational states for optically pumping the aR(6,0) NH3 transition with the 9R16 CO<sub>2</sub> laser line. Recently there have been reports in the literature<sup>2,3</sup> of additional NH3 laser transitions observed at 12.08  $\mu$ m (sP(7,0)) and 12.2  $\mu$ m (aP(6,0)).

We have demonstrated the necessity of simultaneous FIR laser oscillation to produce laser action on the 12.2  $\mu m$  transition. Several optical cavity configurations were used to study the laser process. In the simplest arrangement the  $CO_2$  laser pump is injected into the ammonia cavity through a small hole in the entrance mirror. Ammonia emission is coupled out through this same opening or through a hole in the rear mirror. Both mirrors are mounted internally at the ends of a copper or stainless steel tube.

Figure 2 illustrates the pressure and threshold-time behavior of the 12.2  $\mu m$  and 12.8  $\mu m$  NH3 laser transitions. For these measurements a Lumonics (Model 103) TEA laser with an cutput energy of 250 mJ on 9R16 was used. Figure 2a shows the CO<sub>2</sub> laser pulse shape (200 nsec/div). As shown in Figure 2b, at low pressures (0.3 torr), both the 12.2  $\mu m$  and 12.8  $\mu m$  ammonia laser transitions were present. The intensity ratio (12.2/12.8), of the lines was 0.01. At low pressures, the total pulse duration for the 12.2  $\mu m$  line was 2  $\mu$ sec and 1  $\mu$ sec for relative intensity of the two lines changed to 0.5, and the duration of the 12.8  $\mu$ m line was reduced to 200nsec, while the 12.2  $\mu$ m pulse duration remained about 2  $\mu$ sec. For pressures greater than 1 torr, up to the highest NH3 pressure (30 torr) where laser emission was observed, only the 12.2  $\mu$ m line was observed. Figure 2d shows the 12.2  $\mu$ m pulse shape at 1.3 and 10 torr.

A 12.5 mm thick KCl Brewster angle window was inserted between one of the cavity mirrors and the NH3 laser cavity. It was now found that only the 12.8  $\mu$ m NH3 transition appeared regardless of pressure or buffer gas (He or N<sub>2</sub>) used. KCl does not transmit infrared radiation at wavelengths longer than 30  $\mu$ m. NH3 FIR laser emission wavelengths are in the range of 70 - 150  $\mu$ m for optical rumping with the 9R16 CO<sub>2</sub> line<sup>3,4</sup>. Hence, blocking the intracavity FIR laser emission has led to the elimination of the 12.2  $\mu$ m

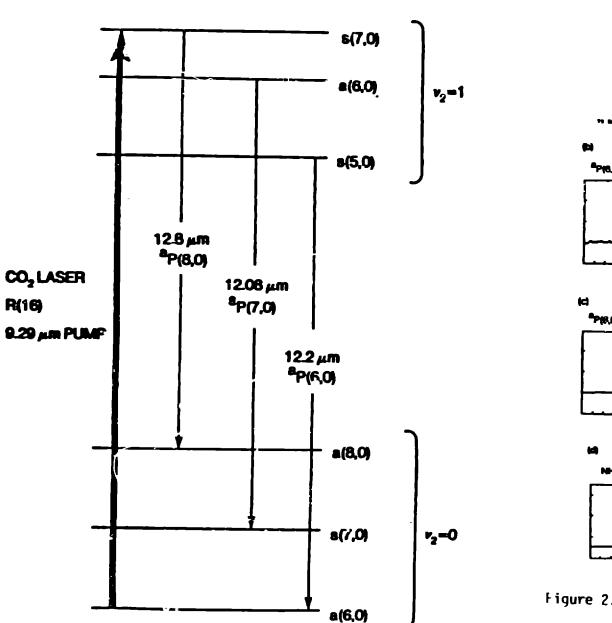


Figure 1 of 4

Figure 1. Relevant energy levels of NH<sub>3</sub> and summary of observed optically pumped mid-infrared laser transitions found in this work and in References 1, 2 and 3.

79 02544

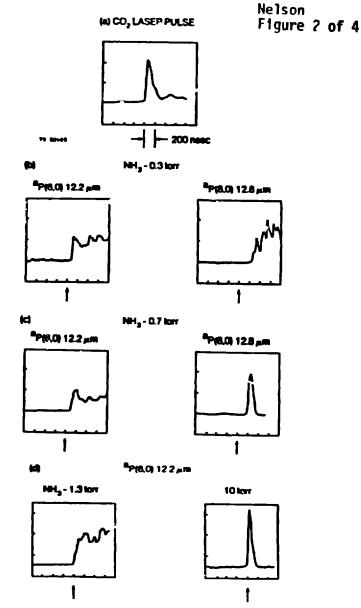
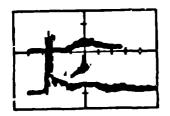
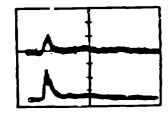


Figure 2. Fulse shapes and relative time histories of the CO<sub>2</sub> pump and NH<sub>3</sub> laser transitions. Onset time of CO<sub>2</sub> laser pulse is indicated by an arrow for each trace. Time scale is 200 nsec/division.



(a) TOP = IGI<sub>3</sub> - 12.2 pm BOTTOM - CO<sub>2</sub> - 9.29 pm 200 nsec/div 1.8 torr IBI<sub>3</sub>



(b) TOP - TOTAL FIR BOTTOM - CO<sub>2</sub> 200 nse:/div

79 82997

Figure 3. Pulse shapes and time histories for the 12.2 m  $\rm IH_3$  laser transition, FIR laser emission and  $\rm CO_2$  pump pulse.

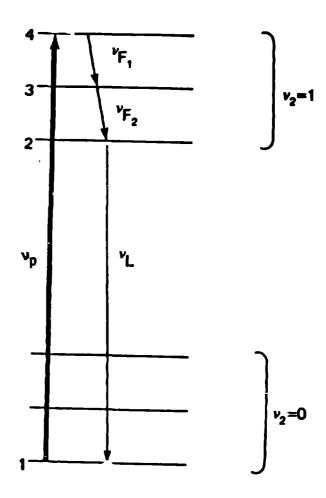


Figure 4. Optically pumped ammonia laser mechanism suggested by present experiments.

mid-infrared laser transition.

In an ammonia laser cavity which used a Littrow grating as a dichroic beamsplitter for the CO2 and ammonia lasers, we have obtained several millijoules of 12.2  $\mu m$  emission. In these studies the 12.2  $\mu m$  emission was not coincident with the gain switched 9R16 CO2 laser pulse. Instead it overlapped the N2-tail portion of the CO2 laser pulse (Figure 3a). FIR emission (Figure 3b), however, followed both the gain switched peak and N2-tail portions of the CO2 laser. This temporal behavior of the 12.2  $\mu m$  emission is not understood at this time and may be related to excessive pump power broadening of the gain  $^5$ .

Figure 4 presents a schematic set of energy levels relevant to the proposed laser mechanism.  $\nu_{\rm p}$  corresponds to the CO<sub>2</sub> pump laser frequency,  $\nu_{\rm L}$  is the mid-infrared NH<sub>3</sub> emission frequency, and the intermediate FIR steps are denoted as  $\nu_{\rm Fl}$  and  $\nu_{\rm F2}$ . The Brewster angle window experiment demonstrates that at least one of the FIR transitions is radiative. Collisional processes may contribute to population transfer into level 2 in Figure 5. A density matrix calculation for the four wave, four level ammonia system indicates that emission at  $\nu_{\rm L}$  will take place near line center and in the absence of a population inversion for levels 1 and 2.

Research support from the U.S. Department of Energy is gratefully acknowledged.

### REFERENCES

- 1. I.Y. Chang and J.D. McGee, Appl Phys. Lett. 28, 526 (1976).
- 2. J.J. Tiee and C. Wittig, J. Appl. Phys. 49, 61 (1978).
- 3. T. Yoshida, N. Yamabayashi, K. Miyazaki and K. Fujisawa, Opt. Comm. 26, 410 (1978).
- 4. K. Gullberg, B. Hartmann and B. Kleman, Phys. Scr. 8, 177 (1973).
- 5. R.L. Panock and R.J. Temkin, IEEE J. Quant. Electron., QE-13, 425 (1977).
- 6. S.J. Hoverson and L.Y. Nelson (unpublished work).